# RESEARCH ARTICLE

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# On Some Double Integrals of $\overline{H}$ -Function of Two Variables and Their Applications

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# **Abstract**

This paper deals with the evaluation of four integrals of  $\overline{H}$ -function of two variables proposed by Singh and Mandia [7] and their applications in deriving double half-range Fourier series for the  $\overline{H}$ -function of two variables. A multiple integral and a multiple half-range Fourier series of the  $\overline{H}$ -function of two variables are derived analogous to the double integral and double half-range Fourier series of the  $\overline{H}$ -function of two variables.

**Key words:**  $\overline{H}$  -function of two variables, Half-range Fourier series,  $\overline{H}$  -function, Multiple half-range Fourier series.

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# I. Introduction

The  $\overline{H}$  -function of two variables will be defined and represented by Singh and Mandia [7] in the following manner:

$$\overline{H}[x,y] = \overline{H}\begin{bmatrix} x \\ y \end{bmatrix} = \overline{H}_{p_{1},q_{1}:p_{2},q_{2}:p_{2},q_{2}}^{o,n_{1}: m_{2},n_{2}:m_{3},n_{2}} \begin{bmatrix} x \\ y \\ (b_{j},\beta_{j};B_{j})_{1,p_{1}},(c_{j},\gamma_{j};K_{j})_{1,n_{2}},(c_{j},\gamma_{j})_{m_{2}+1,p_{2}},(e_{j},E_{j};R_{j})_{1,n_{3}},(e_{j},E_{j})_{m_{3}+1,p_{3}} \end{bmatrix} 
= -\frac{1}{4\pi^{2}} \int_{L_{1}} \int_{L_{2}} \phi_{1}(\xi,\eta) \phi_{2}(\xi) \phi_{3}(\eta) x^{\xi} y^{\eta} d\xi d\eta$$
(1.1)

Where

$$\phi_{1}(\xi,\eta) = \frac{\prod_{j=1}^{n_{1}} \Gamma(1-a_{j}+\alpha_{j}\xi+A_{j}\eta)}{\prod_{j=1}^{p_{1}} \Gamma(a_{j}-\alpha_{j}\xi-A_{j}\eta) \prod_{j=1}^{q_{1}} \Gamma(1-b_{j}+\beta_{j}\xi+B_{j}\eta)}$$

$$(1.2)$$

$$\phi_{2}(\xi) = \frac{\prod_{j=1}^{n_{2}} \left\{ \Gamma\left(1 - c_{j} + \gamma_{j} \xi\right) \right\}^{K_{j}} \prod_{j=1}^{m_{2}} \Gamma\left(d_{j} - \delta_{j} \xi\right)}{\prod_{j=n_{2}+1}^{p_{2}} \Gamma\left(c_{j} - \gamma_{j} \xi\right) \prod_{j=m_{2}+1}^{q_{2}} \left\{ \Gamma\left(1 - d_{j} + \delta_{j} \xi\right) \right\}^{L_{j}}}$$
(1.3)

$$\phi_{3}(\eta) = \frac{\prod_{j=1}^{n_{3}} \left\{ \Gamma\left(1 - e_{j} + E_{j}\eta\right) \right\}^{R_{j}} \prod_{j=1}^{m_{3}} \Gamma\left(f_{j} - F_{j}\eta\right)}{\prod_{j=n_{3}+1}^{p_{3}} \Gamma\left(e_{j} - E_{j}\eta\right) \prod_{j=m_{3}+1}^{q_{3}} \left\{ \Gamma\left(1 - f_{j} + F_{j}\eta\right) \right\}^{S_{j}}}$$
(1.4)

Where x and y are not equal to zero (real or complex), and an empty product is interpreted as unity  $p_i,q_i,n_i,m_j$  are non-negative integers such that  $0 \le n_i \le p_i,o \le m_j \le q_j$  (i=1,2,3;j=2,3). All the  $a_i(j=1,2,...,p_1),b_i(j=1,2,...,q_1),c_i(j=1,2,...,p_2),d_i(j=1,2,...,q_2),$ 

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 $e_{j}(j = 1, 2, ..., p_{3}), f_{j}(j = 1, 2, ..., q_{3})$  are complex parameters.

 $\gamma_j \ge 0$  ( $j = 1, 2, ..., p_2$ ),  $\delta_j \ge 0$  ( $j = 1, 2, ..., q_2$ ) (not all zero simultaneously), similarly

 $E_j \ge 0$  ( $j = 1, 2, ..., p_3$ ),  $F_j \ge 0$  ( $j = 1, 2, ..., q_3$ ) (not all zero simultaneously). The exponents

 $K_{j}(j=1,2,...,n_{3}), L_{j}(j=m_{2}+1,...,q_{2}), R_{j}(j=1,2,...,n_{3}), S_{j}(j=m_{3}+1,...,q_{3})$  can take on non-negative values.

The contour  $L_1$  is in  $\xi$ -plane and runs from  $-i\infty$  to  $+i\infty$ . The poles of  $\Gamma\Big(d_j-\delta_j\xi\Big)(j=1,2,...,m_2)$  lie to the right and the poles of  $\Gamma\Big\{\Big(1-c_j+\gamma_j\xi\Big)\Big\}^{K_j}$   $(j=1,2,...,n_2), \Gamma\Big(1-a_j+\alpha_j\xi+A_j\eta\Big)(j=1,2,...,n_1)$  to the left of the contour. For  $K_j$   $(j=1,2,...,n_2)$  not an integer, the poles of gamma functions of the numerator in (1.3) are converted to the branch points.

The contour  $L_2$  is in  $\eta$ -plane and runs from  $-i\infty$  to  $+i\infty$ . The poles of  $\Gamma(f_j-F_j\eta)(j=1,2,...,m_3)$  lie to the right and the poles of

$$\Gamma\left\{\left(1-e_{j}+E_{j}\eta\right)\right\}^{R_{j}}\ (j=1,2,...,n_{3}), \\ \Gamma\left(1-a_{j}+\alpha_{j}\xi+A_{j}\eta\right)(j=1,2,...,n_{1}) \ \ \text{to the left of the contour.}$$

For  $R_j$  ( $j = 1, 2, ..., n_3$ ) not an integer, the poles of gamma functions of the numerator in (1.4) are converted to the branch points.

The functions defined in (1.1) is an analytic function of x and y, if

$$U = \sum_{j=1}^{p_1} \alpha_j + \sum_{j=1}^{p_2} \gamma_j - \sum_{j=1}^{q_1} \beta_j - \sum_{j=1}^{q_2} \delta_j < 0$$
 (1.5)

$$V = \sum_{j=1}^{p_1} A_j + \sum_{j=1}^{p_3} E_j - \sum_{j=1}^{q_1} B_j - \sum_{j=1}^{q_3} F_j < 0$$
(1.6)

The integral in (1.1) converges under the following set of conditions:

$$\Omega = \sum_{j=1}^{n_1} \alpha_j - \sum_{j=n_1+1}^{p_1} \alpha_j + \sum_{j=1}^{m_2} \delta_j - \sum_{j=m_2+1}^{q_2} \delta_j L_j + \sum_{j=1}^{n_2} \gamma_j K_j - \sum_{j=n_2+1}^{p_2} \gamma_j - \sum_{j=1}^{q_1} \beta_j > 0$$
(1.7)

$$\Lambda = \sum_{i=1}^{n_1} A_j - \sum_{i=n,+1}^{p_1} A_j + \sum_{i=1}^{m_2} F_j - \sum_{i=m,+1}^{q_2} F_j S_j + \sum_{i=1}^{n_3} E_j R_j - \sum_{i=n,+1}^{p_3} E_j - \sum_{i=1}^{q_1} B_j > 0$$
(1.8)

$$|\arg x| < \frac{1}{2}\Omega\pi, |\arg y| < \frac{1}{2}\Lambda\pi$$
 (1.9)

The behavior of the  $\overline{H}$  -function of two variables for small values of |z| follows as:

$$\overline{H}[x, y] = 0(|x|^{\alpha}|y|^{\beta}), \max\{|x|, |y|\} \to 0$$
 (1.10)

Where

$$\alpha = \min_{1 \le j \le m_2} \left[ \operatorname{Re} \left( \frac{d_j}{\delta_j} \right) \right] \qquad \beta = \min_{1 \le j \le m_2} \left[ \operatorname{Re} \left( \frac{f_j}{F_j} \right) \right]$$
(1.11)

For large value of |z|,

$$\overline{H}[x, y] = 0\{|x|^{\alpha'}, |y|^{\beta'}\}, \min\{|x|, |y|\} \to 0$$
(1.12)

Where

$$\alpha' = \max_{1 \le j \le n_2} \operatorname{Re} \left( K_j \frac{c_j - 1}{\gamma_j} \right), \ \beta' = \max_{1 \le j \le n_3} \operatorname{Re} \left( R_j \frac{e_j - 1}{E_j} \right)$$
(1.13)

Provided that U < 0 and V < 0.

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If we take

$$K_{j} = 1(j = 1, 2, ..., n_{2}), L_{j} = 1(j = m_{2} + 1, ..., q_{2}), R_{j} = 1(j = 1, 2, ..., n_{3}), S_{j} = 1(j = m_{3} + 1, ..., q_{3})$$
 in

(1.1), the  ${\cal H}$  -function of two variables reduces to  ${\cal H}$  -function of two variables due to [4]. Orthogonality of Sine and Cosine functions:

$$\int_{0}^{\pi} \sin mx \sin nx dx = \frac{0; m \neq n}{\frac{\pi}{2}; m = n}$$
 (1.14)

$$\int_{0}^{\pi} \cos mx \cos nx dx = \frac{0; m \neq n}{\frac{\pi}{2}; m = n}$$

$$\frac{\pi}{\pi} = \frac{1}{2} = \frac{\pi}{m = n = 0}$$
(1.15)

#### II. Double Integrals

The following double integrals has been evaluated in this paper

$$\int_{0}^{\pi} \int_{0}^{\pi} (\sin x)^{\lambda - 1} (\sin y)^{\mu - 1} \sin rx \sin ty \, h(x, y) \, dx dy$$

$$=2^{2-\lambda-\mu}\pi^{2}\sin\left(\frac{r\pi}{2}\right)\sin\left(\frac{t\pi}{2}\right)\psi(r,t)$$
(2.1)

$$\int_{0}^{\pi} \int_{0}^{\pi} (\sin x)^{\lambda - 1} (\sin y)^{\mu - 1} \sin rx \cos ty \, h(x, y) \, dx dy$$

$$=2^{2-\lambda-\mu}\pi^2\sin\left(\frac{r\pi}{2}\right)\cos\left(\frac{t\pi}{2}\right)\psi(r,t)$$
(2.2)

$$\int_{0}^{\pi} \int_{0}^{\pi} (\sin x)^{\lambda - 1} (\sin y)^{\mu - 1} \cos rx \sin ty \, h(x, y) \, dx dy$$

$$=2^{2-\lambda-\mu}\pi^2\cos\left(\frac{r\pi}{2}\right)\sin\left(\frac{t\pi}{2}\right)\psi(r,t)$$
(2.3)

$$\int_{0}^{\pi} \int_{0}^{\pi} (\sin x)^{\lambda - 1} (\sin y)^{\mu - 1} \cos rx \cos ty \, h(x, y) \, dx dy$$

$$=2^{2-\lambda-\mu}\pi^2\cos\left(\frac{r\pi}{2}\right)\cos\left(\frac{t\pi}{2}\right)\psi(r,t)$$
(2.4)

Where

$$h(x,y) = \overline{H}_{p_1,q_1:p_2,q_2:p_2,q_2}^{o,n_1: m_2,n_2:m_3,n_2} \begin{bmatrix} u(\sin x)^{2c} & (a_j,\alpha_j;A_j)_{1,p_1}, (c_j,\gamma_j;K_j)_{1,p_2}, (c_j,\gamma_j)_{n_2+1,p_2}, (e_j,E_j;R_j)_{1,n_3}, (e_j,E_j)_{n_3+1,p_3} \\ (b_j,\beta_j;B_j)_{1,q_1}, (d_j,\delta_j)_{1,m_2}, (d_j,\delta_j;L_j)_{m_2+1,q_2}, (f_j,F_j)_{1,m_3}, (f_j,F_j;S_j)_{m_3+1,q_3} \end{bmatrix}$$

And

$$\psi(r,t) =$$

$$\overline{H}_{p_{1}+2,q_{1}+4:p_{2},q_{2};p_{2},q_{2}}^{o,n_{1}+2:} \left[ u2^{2(c+d)} \\ v \\ \left| (1-\lambda,2c;1),(1-\mu,2d;1),\left(a_{j},\alpha_{j};A_{j}\right)_{1,p_{1}},\left(c_{j},\gamma_{j};K_{j}\right)_{1,n_{2}},\left(c_{j},\gamma_{j}\right)_{n_{2}+1,p_{2}},\left(e_{j},E_{j};R_{j}\right)_{1,n_{3}},\left(e_{j},E_{j}\right)_{n_{3}+1,p_{3}} \\ \left(b_{j},\beta_{j};B_{j}\right)_{1,q_{1}},\left(\frac{1-\lambda\pm r}{2},c;1\right),\left(\frac{1-\mu\pm t}{2},d;1\right),\left(d_{j},\delta_{j}\right)_{1,m_{2}},\left(d_{j},\delta_{j};L_{j}\right)_{m_{2}+1,q_{2}},\left(f_{j},F_{j}\right)_{1,m_{3}},\left(f_{j},F_{j};S_{j}\right)_{m_{3}+1,q_{3}} \\ \right]$$

And  $(\lambda \pm \mu)$  stands for the pair of parameters  $(\lambda + \mu), (\lambda - \mu)$ .

Also 
$$\operatorname{Re}(\lambda) + 2c \min_{1 \leq j \leq m_1} \left( \frac{b_j}{\beta_j} \right) + 2c \min_{1 \leq j \leq m_2} \left( \frac{d_j}{\delta_j} \right) > 0$$
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$$\operatorname{Re}(\mu) + 2d \min_{1 \leq j \leq m_1} \left( \frac{b_j}{\beta_j} \right) + 2c \min_{1 \leq j \leq m_2} \left( \frac{d_j}{\delta_j} \right) > 0 \text{ and conditions (1.7), (1.8) and (1.9) are also satisfied.}$$

**Proof:** If we express the H -function of two variables occurring in the integrand of (2.1) as the Mellin-Barnes type integral (1.1) and interchange the order of integrations (which is permissible due to the absolute convergence of the integrals involved in the process), we get L.H.S. of (2.1)

$$= -\frac{1}{4\pi^{2}} \int_{L_{1}} \int_{L_{2}} \phi_{1}(\xi, \eta) \phi_{2}(\xi) \phi_{3}(\eta) u^{\xi} v^{\eta}$$

$$\left[ \int_{0}^{\pi} (\sin x)^{\mu + 2c\xi - 1} \sin rx dx \int_{0}^{\pi} (\sin y)^{\mu + 2d\xi - 1} \sin ty dy \right] d\xi d\eta$$

Now, applying the result ([5],p.70, eq. (3.1.5)) and equation (1.1), the result (2.1) follows at once. The remaining integrals can be evaluated similarly.

# III. Double Half-Range Fourier Series

The following double half-range Fourier series will be proved:

$$f(x,y) = 2^{4-\lambda-\mu} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \sin\left(\frac{m\pi}{2}\right) \sin\left(\frac{n\pi}{2}\right) \psi(m,n) \sin mx \sin ny$$
(3.1)

$$f(x,y) = 2^{4-\lambda-\mu} \sum_{m=1}^{\infty} \sum_{n=0}^{\infty} \sin\left(\frac{m\pi}{2}\right) \cos\left(\frac{n\pi}{2}\right) \psi(m,n) \sin mx \cos ny$$
(3.2)

$$f(x,y) = 2^{4-\lambda-\mu} \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} \cos\left(\frac{m\pi}{2}\right) \sin\left(\frac{n\pi}{2}\right) \psi(m,n) \cos mx \sin ny$$
(3.3)

$$f(x,y) = 2^{4-\lambda-\mu} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \cos\left(\frac{m\pi}{2}\right) \cos\left(\frac{n\pi}{2}\right) \psi(m,n) \cos mx \cos ny$$
(3.4)

Where  $f(x, y) = (\sin x)^{\lambda - 1} (\sin y)^{\mu - 1} h(x, y)$  and provided that

$$\operatorname{Re}(\lambda) + 2c \min_{1 \leq j \leq m_1} \left( \frac{b_j}{\beta_j} \right) + 2c \min_{1 \leq j \leq m_2} \left( \frac{d_j}{\delta_j} \right) > 0 \,, \, \operatorname{Re}(\mu) + 2d \min_{1 \leq j \leq m_1} \left( \frac{b_j}{\beta_j} \right) + 2c \min_{1 \leq j \leq m_2} \left( \frac{d_j}{\delta_j} \right) > 0 \, \text{and}$$

conditions (1.7), (1.8) and (1.9) are also satisfied.

**Proof:** To prove (3.1), let

$$f(x,y) = (\sin x)^{\lambda-1} (\sin y)^{\mu-1} h(x,y) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{m,n} \sin mx \sin ny$$
 (3.5)

Which is valid since f(x, y) is continuous and of bounded variation in the open interval  $(0, \pi)$ .

Multiplying both sides of (3.5) by  $\sin rx \sin ty$  and integrating from 0 to  $\pi$  with respect to both x and y, it is seen that

$$\int_{0}^{\pi} \int_{0}^{\pi} (\sin x)^{\lambda - 1} (\sin y)^{\mu - 1} h(x, y) \sin rx \sin ty dx dy =$$

$$\sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{m,n} \int_{0}^{\pi} \int_{0}^{\pi} \sin mx \sin ny \sin rx \sin ty dx dy$$
(3.6)

Now using (2.1), and orthogonal property of sine functions, it follows that

$$A_{r,t} = 2^{4-\lambda-\mu} \sin\left(\frac{r\pi}{2}\right) \sin\left(\frac{t\pi}{2}\right) \psi(r,t)$$
(3.7)

Substituting the value of  $A_{m,n}$  from (3.7) to (3.5), the result (3.1) follows at once.

To establish (3.2), put

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$$f(x,y) = \sum_{m=1}^{\infty} \sum_{n=0}^{\infty} B_{m,n} \sin mx \cos ny$$
 (3.8)

Multiplying both sides of (3.8) by  $\sin rx \cos ty$  and integrating from 0 to  $\pi$  with respect to both x and y and using (2.2) and orthogonal properties of sine and cosine functions, we find that

$$B_{r,t} = 2^{4-\lambda-\mu} \sin\left(\frac{r\pi}{2}\right) \cos\left(\frac{t\pi}{2}\right) \psi(r,t)$$
(3.9)

Except that  $B_{r,0}$  is one-half of the above value.

From (3.8) and (3.9), the series (3.2) follows easily. The result (3.3) can be established in a manner similar to above.

To establish (3.4), let

$$f(x,y) = \sum_{m=1}^{\infty} \sum_{n=0}^{\infty} D_{m,n} \cos mx \cos ny$$
 (3.10)

Multiplying both sides of (3.8) by  $\cos rx \cos ty$  and integrating from 0 to  $\pi$  with respect to both x and y and using (2.2) and orthogonal properties of cosine functions, we obtain

$$D_{r,t} = 2^{4-\lambda-\mu} \cos\left(\frac{r\pi}{2}\right) \cos\left(\frac{t\pi}{2}\right) \psi(r,t)$$
(3.11)

Except that  $D_{0,t}, D_{r,0}$  are one-half and  $D_{0,0}$  is quarter of the above values.

The Fourier series (3.4) now follows from (3.10) and (3.11).

# **IV. Multiple Integrals**

The following multiple integral analogous to (2.1) can be derived easily on following the procedure as given \$2 and taking the help of ([5],p.70,eq.(3.1.5):

$$\int_{0}^{\pi} ... \int_{0}^{\pi} (\sin x_{1})^{\lambda_{1}-1} (\sin x_{n})^{\lambda_{n}-1} \sin r_{1} x_{1} ... \sin r_{n} x_{n} h(x_{1},...,x_{n}) dx_{1} ... dx_{n}$$

$$=2^{n-\lambda_1-...-\lambda_n}\pi^n\sin\left(\frac{r_1\pi}{2}\right)...\sin\left(\frac{r_n\pi}{2}\right)\psi(r_1,...,r_n)$$
(4.1)

Where 
$$\operatorname{Re}(\lambda_i) + 2c \min_{1 \le j \le m_1} \left( \frac{b_j}{\beta_i} \right) + 2c \min_{1 \le j \le m_2} \left( \frac{d_j}{\delta_j} \right) > 0$$
;  $i = 1, 2, ..., n$  and conditions (1.7), (1.8) and (1.9) are

also hold, and

$$h(x_1,...,x_n) = \overline{H}_{p_1,q_1:p_2,q_2:p_2,q_2}^{o,n_1: m_2,n_2:m_3,n_2} \begin{bmatrix} uv(\sin x_1)^{2c_1} \\ \vdots \\ (\sin x_n)^{2c_n} \end{bmatrix} \begin{bmatrix} (a_j,\alpha_j;A_j)_{1,p_1}, (c_j,\gamma_j;K_j)_{1,n_2}, (c_j,\gamma_j)_{n_2+1,p_2}, (e_j,E_j;R_j)_{1,n_3}, (e_j,E_j)_{n_3+1,p_3} \\ (b_j,\beta_j;B_j)_{1,q_1}, (d_j,\delta_j)_{1,m_2}, (d_j,\delta_j;L_j)_{m_2+1,q_2}, (f_j,F_j)_{1,m_3}, (f_j,F_j;S_j)_{m_3+1,q_3} \end{bmatrix}$$

And

$$\psi(r_1,...,r_n) =$$

$$\overline{H}{}^{o,n_1+2:}{}^{m_2,n_2:m_3,n_2}_{p_1+2:q_1+4:p_2,q_2;p_2,q_2} \begin{bmatrix} u2^{2(c_1+\ldots+c_n)} \\ v \\ (b_j,\beta_j;B_j)_{1,q_1}, \underbrace{\left(1-\lambda_j,2c_j;1\right)_{1,n},\left(a_j,\alpha_j;A_j\right)_{1,p_1},\left(c_j,\gamma_j;K_j\right)_{1,n_2},\left(c_j,\gamma_j\right)_{n_2+1,p_2},\left(e_j,E_j;R_j\right)_{1,n_3},\left(e_j,E_j\right)_{n_3+1,p_3}} \\ \left(b_j,\beta_j;B_j\right)_{1,q_1},\underbrace{\left(1-\lambda_j\pm r_j,c_j;1\right)_{1,n_2},\left(d_j,\delta_j\right)_{1,m_2},\left(d_j,\delta_j;L_j\right)_{m_2+1,q_2},\left(f_j,F_j\right)_{1,m_3},\left(f_j,F_j;S_j\right)_{m_3+1,q_3}} \right]$$

For j = 1, 2, ..., n.

Multiple integrals analogous to (2.2) to (2.4) can also be written easily.

### V. Multiple Half-Range Fourier Series

The following multiple half-range Fourier series analogous to (2.1) can be derived on following lines as given in \$3 [5], using the integral (4.1) and the multiple orthogonal property of sine functions:

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$$f(x_1,...,x_n) = 2^{2n-\lambda_1-...-\lambda_n} \sum_{m_1=1}^{\infty} ... \sum_{m_n=1}^{\infty} \sin\left(\frac{m_1\pi}{2}\right) ... \sin\left(\frac{m_n\pi}{2}\right)$$

$$\psi(m_1,...,m_n)\sin(m_1x_1)...\sin(m_nx_n)$$

Where 
$$\operatorname{Re}(\lambda_i) + 2c \min_{1 \leq j \leq m_1} \left( \frac{b_j}{\beta_j} \right) + 2c \min_{1 \leq j \leq m_2} \left( \frac{d_j}{\delta_j} \right) > 0$$
;  $i = 1, 2, ..., n$  and conditions (1.7), (1.8) and (1.9) are

also satisfied, and  $f(x_1,...,x_n) = (\sin x_1)^{\lambda_1-1}...(\sin x_n)^{\lambda_n-1}h(x_1,...,x_n)$  Similarly the multiple half-range Fourier series analogous to (2.2), (2.3) and (2.4) can also be solved.

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